



Instituto Nacional de Pesquisas Espaciais

Parametrização de chuva e nuvens

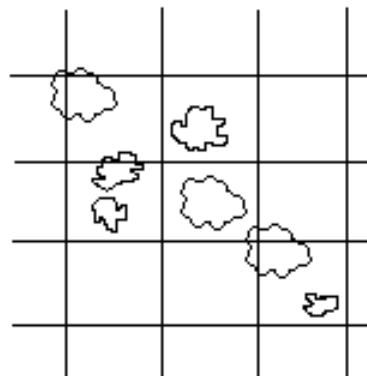
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**Entrenamiento em Modelado de Escenarios de Cambio Climático
Cachoeira Paulista, 30 de agosto a 4 de setembro de 2009**

The need for cumulus parameterization

- *Convective clouds can organise in clusters and show their collective effects in model grid-box.



- *Large-scale destabilizes the environment >>> cumulus parameterization scheme acts to remove the convective instability

- *The upward fluxes of heat, moisture and momentum in the cloud can be seen by means of an area average over the equations of mass continuity and heat energy.

- *Up to which resolution should the parameterization act in a model? No clear agreement.

Médias

- Os movimentos atmosféricos existem em várias escalas espaciais. Em um modelo numérico há processos resolvidos pela grade do modelo e outros processos “sub-grade”.
- Há necessidade de descrever os processos **resolvidos** pelo sistema de observações e aqueles **não-resolvidos** e designados por perturbação (“eddy”).
- Para identificar as propriedades estatísticas de um sistema, utiliza-se médias.

•1 Média no volume da grade (Grid-volume averaging):

$$\phi = \bar{\phi} + \phi'$$

Variável composta por uma média resolvida e uma perturbação

‘ $\bar{}$ ’ se refere à média na grade do modelo

ϕ' perturbação subgrade

$$\bar{\phi} = \int_t^{t+\Delta t} \int_x^{x+\Delta x} \int_y^{y+\Delta y} \int_z^{z+\Delta z} \phi dz dy dx dt / (\Delta t \Delta x \Delta y \Delta z)$$

$\Delta x, \Delta y, \Delta z$ são as dimensões da grade do modelo e Δt o passo de tempo.

$$\overline{\bar{\phi}} = \bar{\phi}$$

$$\overline{\phi'} = 0$$

$$\overline{\frac{\partial \bar{u}}{\partial t}} = \frac{\partial \bar{u}}{\partial t}$$

$$\overline{\frac{\partial \bar{u}}{\partial x}} = \frac{\partial \bar{u}}{\partial x}$$

$$\overline{\phi' \bar{u}} = 0$$

$$\overline{\phi' w'} \neq 0$$

$$\overline{T' w'} \neq 0$$

A média do produto da correlação das perturbações resulta em valor diferente de zero!!

Equations of heat, moisture and continuity

$$\frac{\partial \bar{\theta}}{\partial t} + \overline{\nabla \cdot \theta \mathbf{v}} + \frac{\partial \bar{\theta} \bar{w}}{\partial z} = \frac{Q_R}{c_p \pi} + \frac{L}{c_p \pi} (c - e) - \frac{1}{\rho} \frac{\partial}{\partial z} \overline{\rho \theta w}$$

$$\frac{\partial \bar{q}}{\partial t} + \overline{\nabla \cdot q \mathbf{v}} + \frac{\partial \bar{q} \bar{w}}{\partial z} = -(c - e) - \frac{1}{\rho} \frac{\partial}{\partial z} \overline{\rho q w}$$

$$\overline{\nabla \cdot \mathbf{v}} + \frac{\partial \bar{w}}{\partial z} = 0$$

These subgrid terms need to be parameterized because their effects contribute to model grid scale

The vertical eddy fluxes are due mainly to the cumulus convection and turbulent motions in the boundary layer.

Effects of cumulus convection on large scale thermodynamic fields

$$Q_1 \equiv \frac{\partial \bar{s}}{\partial t} + \bar{v} \bullet \nabla \bar{s} + \bar{\omega} \frac{\partial \bar{s}}{\partial p} = Q_R + L(\bar{c} - \bar{e}) - \nabla \cdot \overline{s'v'} - \frac{\partial}{\partial p} \overline{s'\omega'}$$

Q_1 : Apparent Heat source

$$Q_2 \equiv -L \left(\frac{\partial \bar{q}}{\partial t} + \bar{v} \bullet \nabla \bar{q} + \bar{\omega} \frac{\partial \bar{q}}{\partial p} \right) = L(\bar{c} - \bar{e}) + L \nabla \cdot \overline{q'v'} + L \frac{\partial}{\partial p} \overline{q'\omega'}$$

Q_2 : Apparent Moisture sink

Assume horizontal fluxes are negligible compared to vertical fluxes

$$Q_1 - Q_R - Q_2 = - \frac{\partial}{\partial p} \overline{h'\omega'}$$

Eddy (sensible and latent) heat fluxes
 h is the moist static energy

$$\frac{c_p}{g} \int_{p_t}^{p_s} (Q_1 - Q_R) dp = LP + H_s$$

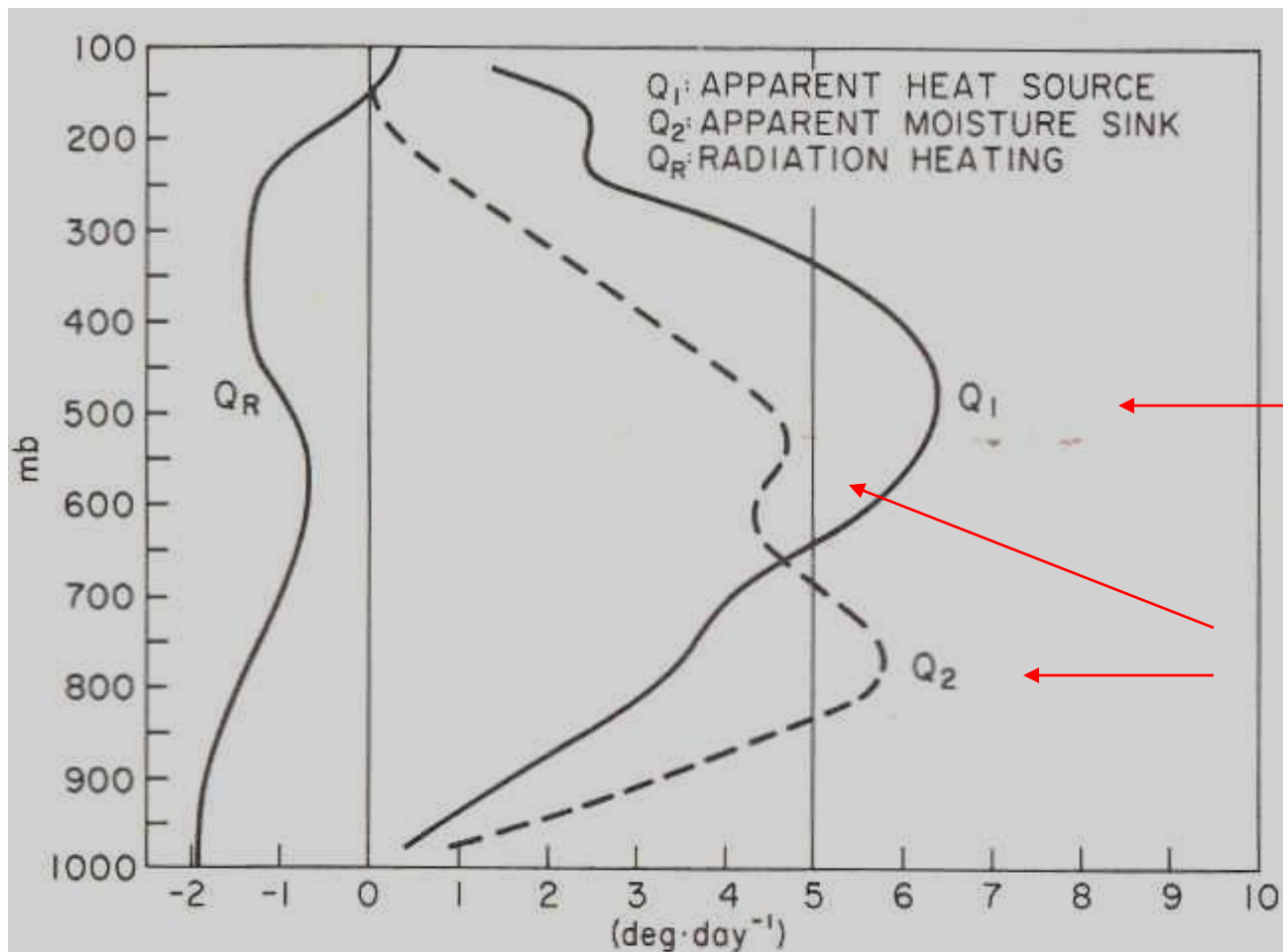
Verify against sfc observations

$$\frac{c_p}{g} \int_{p_t}^{p_s} Q_2 dp = L(P - E)$$

Yanai et al., 1973

Yanai and Johnson, 1993

Fig 4.1, Fig 4.3, F4.17



~5°C/dia,
~400-500 hPa

~4°C/dia,
~500 hPa,
~800 hPa

FIG. 4.1. The mean apparent heat source Q_1 (solid) and moisture sink Q_2 (dashed) over the Marshall Islands. On the left is the radiation heating rate given by Dopplick (1972) (from Yanai et al. 1973).

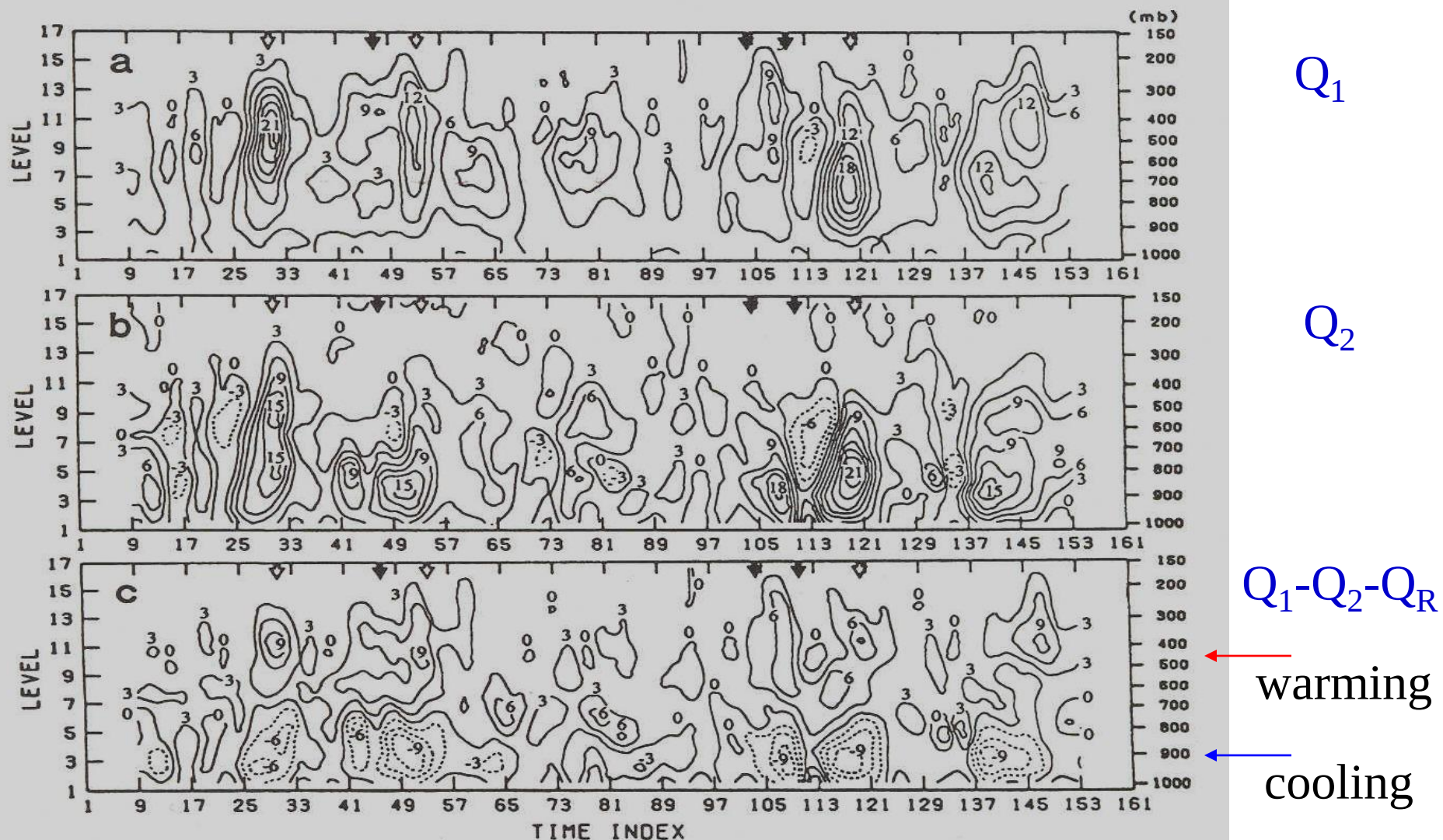


FIG. 4.3. Time-height sections of the observed (a) $Q_1 - Q_R$, (b) Q_2 , and (c) $Q_1 - Q_2 - Q_R$, averaged over the $3^\circ \times 3^\circ$ area at the center of the GATE network during a period from 0000 UTC 31 August (index 9) to 0000 UTC 18 September 1974 (index 153). Units are kelvins per day (from Cheng and Yanai 1989).

$$Q_1 - Q_R - Q_2 = -\frac{\partial}{\partial p} \overline{h' \omega'}$$

TRADEWIND CUMULUS

Downward motion

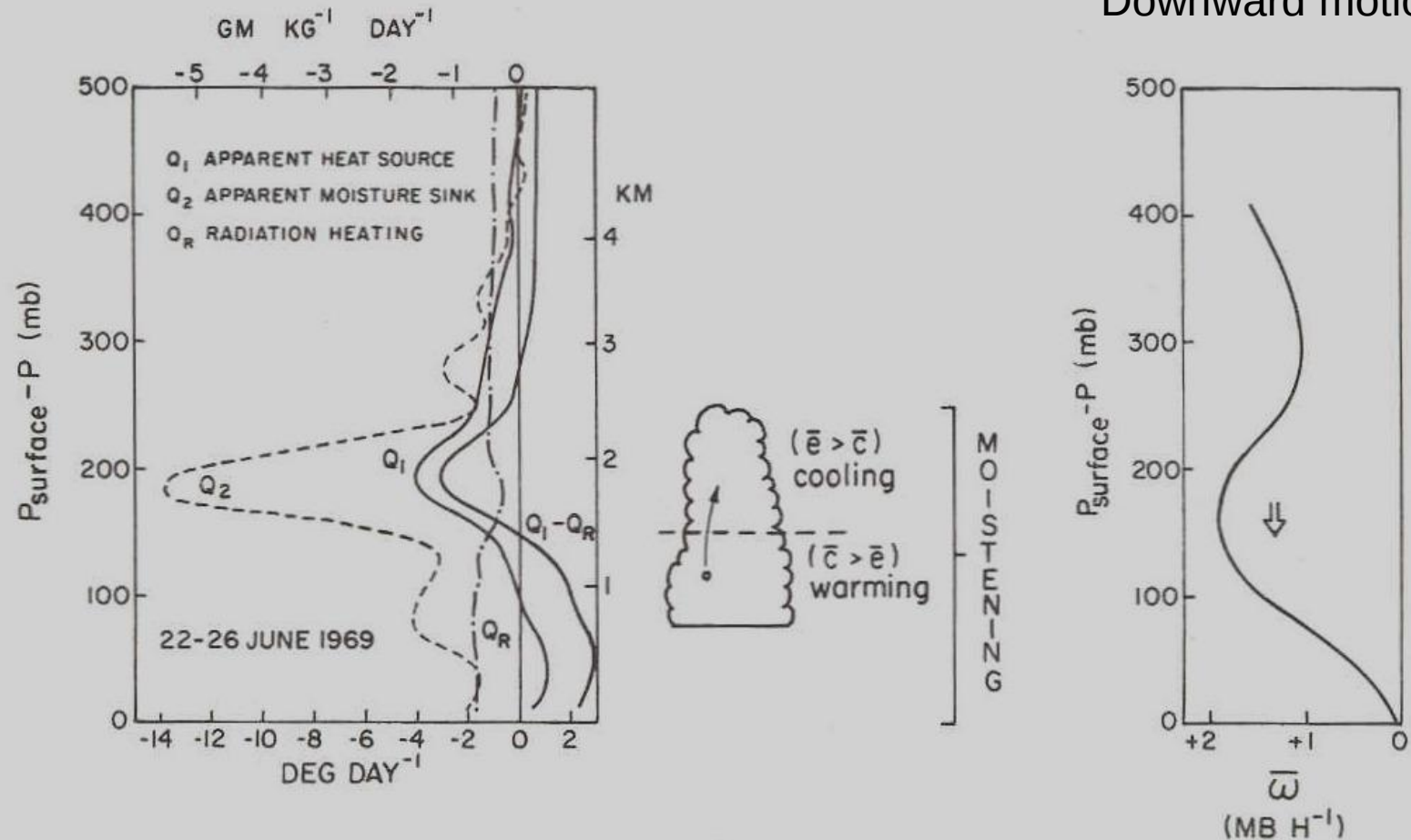


FIG. 4.17. (Left) The observed Q_1 , Q_2 , Q_R , and $Q_1 - Q_R$ for the undisturbed BOMEX period 22-26 June 1969 (from Nitta and Esbensen 1974). (Center) Schematic of trade-wind cumulus layer showing effects of condensation and evaporation on the heat and moisture budgets. (Right) Mean vertical p velocity $\bar{\omega}$ over budget area.

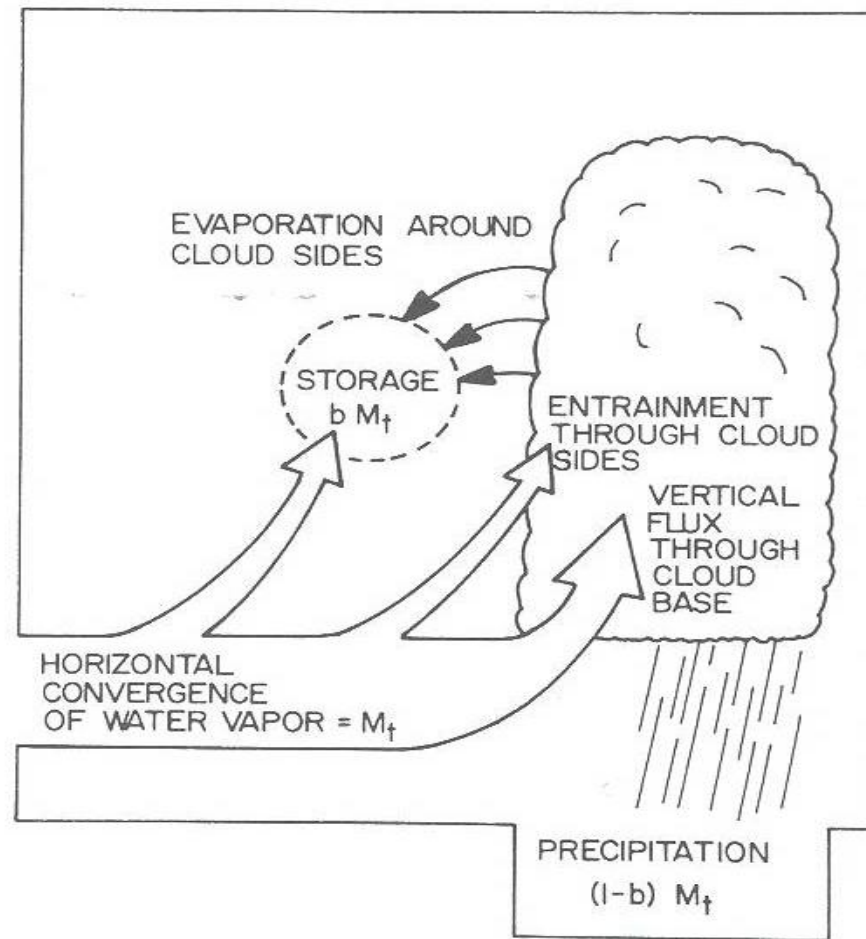
Types of convective scheme:

Adjustment: Betts-Miller (1986), Janjic (1994)

Kuo: Kuo (1974)

Mass-flux: Arakawa-Schubert (1974), Fritsch-Chappell (1980),
Tiedtke (1989), Grell, Kain-Fritsch (1993)

The Kuo Convective Parameterization Scheme



Betts-Miller-Janjic

- The Betts-Miller scheme (Betts and Miller, 1986) uses reference profiles of T and q to relax the model profiles in convective unstable conditions. Profiles derived from campaigns GATE, VIMEX, etc
- The reference T and q profiles are based on observational studies of convective equilibrium in the tropics.
- Treats deep and shallow convection.
- (Modification by Janjic, 1994)

1. Determine type of cloud

- Parcel lift: determine cloud base and cloud top
- Check: cloud depth > 290 hPa: deep convection, else shallow convection.

2. Determine reference profiles

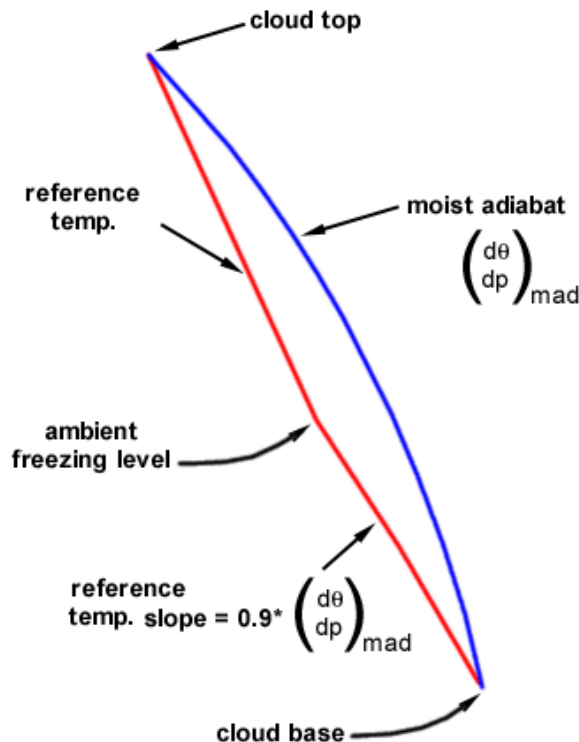
$$\int_{base}^{top} (c_p \Delta T - L \Delta q) = 0$$

Make sure enthalpy is conserved

Deep Convection

Temperature Reference Profile

Construction of
1st Guess Temperature Reference Profile
for Deep Convection



To draw the temperature profile
3 levels are important:

- cloud base
- freezing level
- cloud top

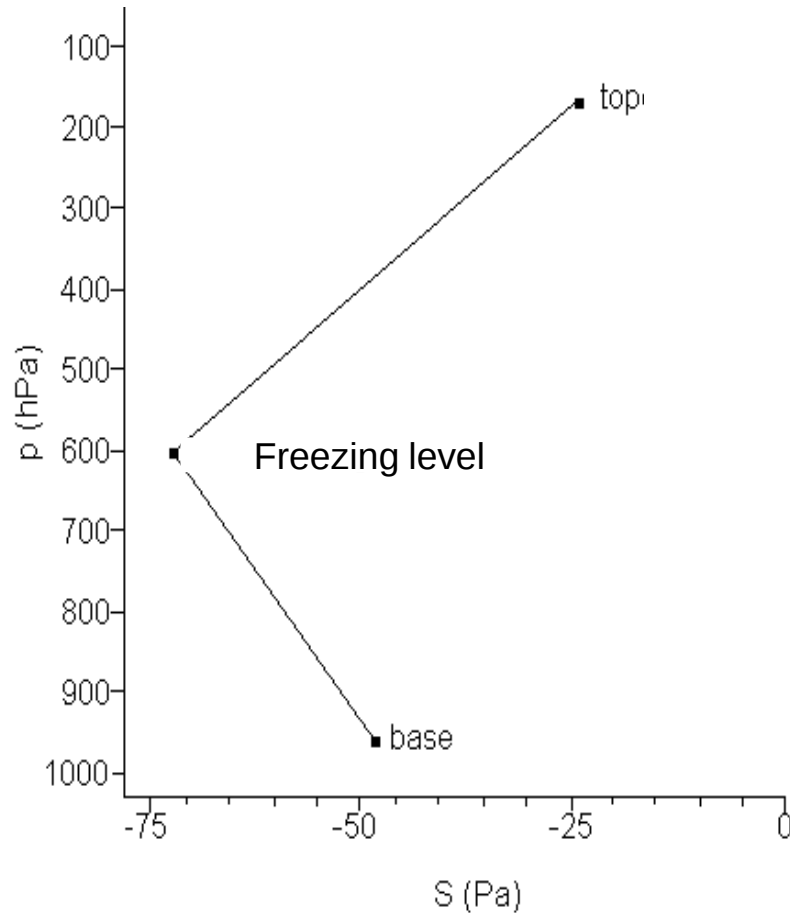
Some instability is left in the
lower part of the cloud.

Profile is linearly interpolated
from the freezing level to the
cloud top

Mike Staudenmaier

Deep Convection

Moisture Reference Profile



$$DSP = p_{sat} - p$$

Deficit saturation pressure (DSP) is defined at the 3 levels.

- cloud base: DSPb
- freezing level: DSP0
- cloud top: DSPt

Values are linearly interpolated between the levels

$$T_{new} = T_{old} + \frac{\Delta t_{cnv}}{\tau} [T_{ref} - T]$$

$$q_{new} = q_{old} + \frac{\Delta t_{cnv}}{\tau} [q_{ref} - q]$$

$$\Delta t_{cnv} = 4 * \Delta t$$

$$\tau = 3000s$$

$$P = \frac{1}{\rho_w g} \frac{\Delta t_{cnv}}{\tau} \sum_{base}^{top} (q_{ref} - q)(p_s - p_t)$$

1. Cloud Efficiency

- Efficiency related to the precipitation production
- Proportional to the cloud column "entropy" change, per unit precipitation produced.
- Efficiency varies from 0.2 to 1.0
- Modifies reference profiles
- Modifies relaxation time

$$\tau' = \frac{\tau}{F(E)}$$

F(E) is linear

$$0.7 \leq F \leq 1.0 \text{ for } 0.2 \leq E \leq 1.0$$

Thus,

larger $E \Rightarrow$ less mature system

smaller $E \Rightarrow$ more mature system

USE E TO MODERATE HEAVY RAIN IN LONG-LIVED MATURE SYSTEMS

(A) Modify the humidity reference profile

(B) Modify the relaxation time τ

$$T_{new} = T_{old} + \frac{\Delta t_{cnv}}{\tau} \left[T_{ref} - T_{old} \right] F(E)$$

$$q_{new} = q_{old} + \frac{\Delta t_{cnv}}{\tau} \left[q_{ref} - q_{old} \right] F(E)$$

OR

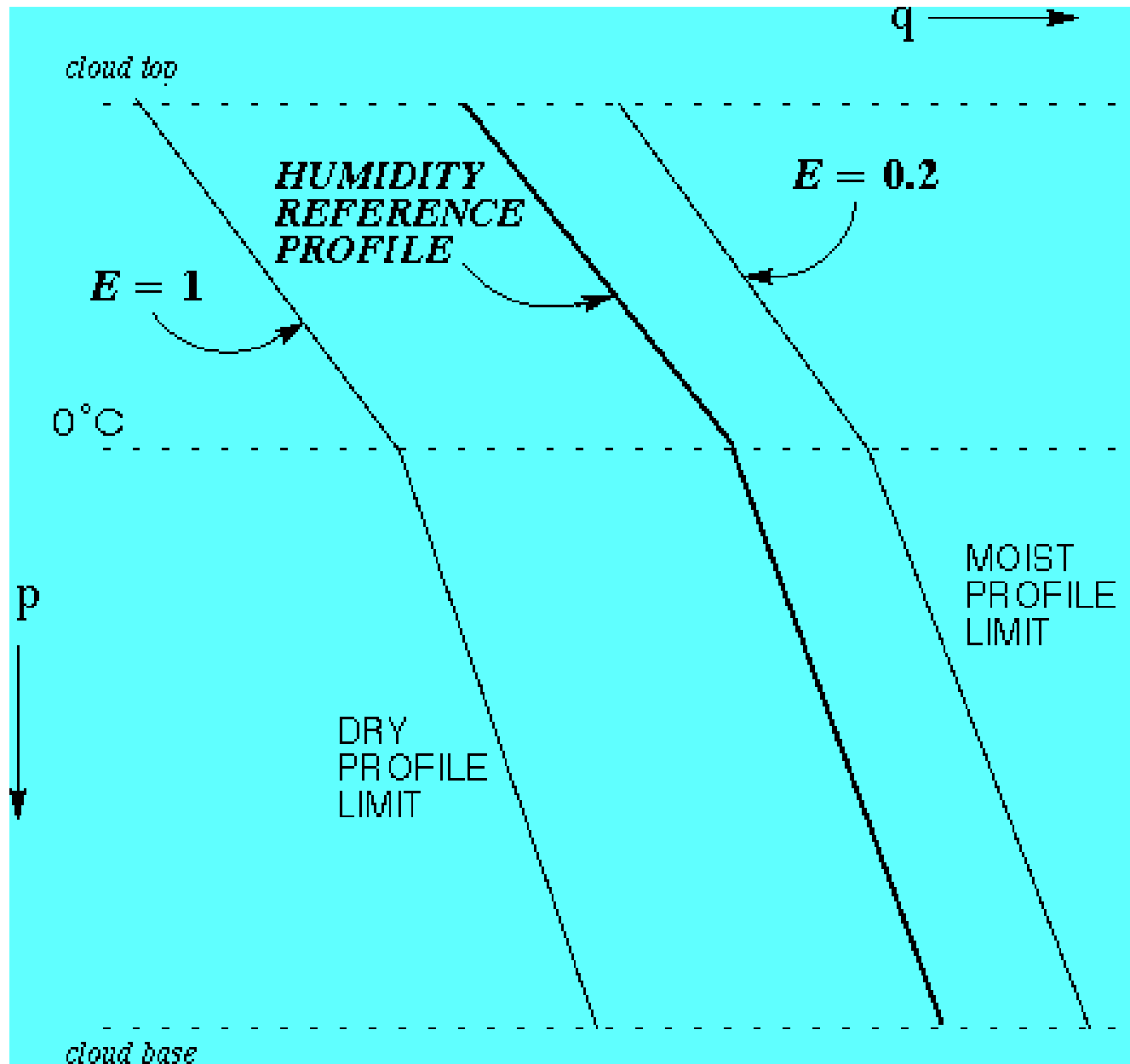
$$T_{new} = T_{old} + \frac{\Delta t_{cnv}}{\tau'} \left[T_{ref} - T_{old} \right]$$

$$q_{new} = q_{old} + \frac{\Delta t_{cnv}}{\tau'} \left[q_{ref} - q_{old} \right]$$

where

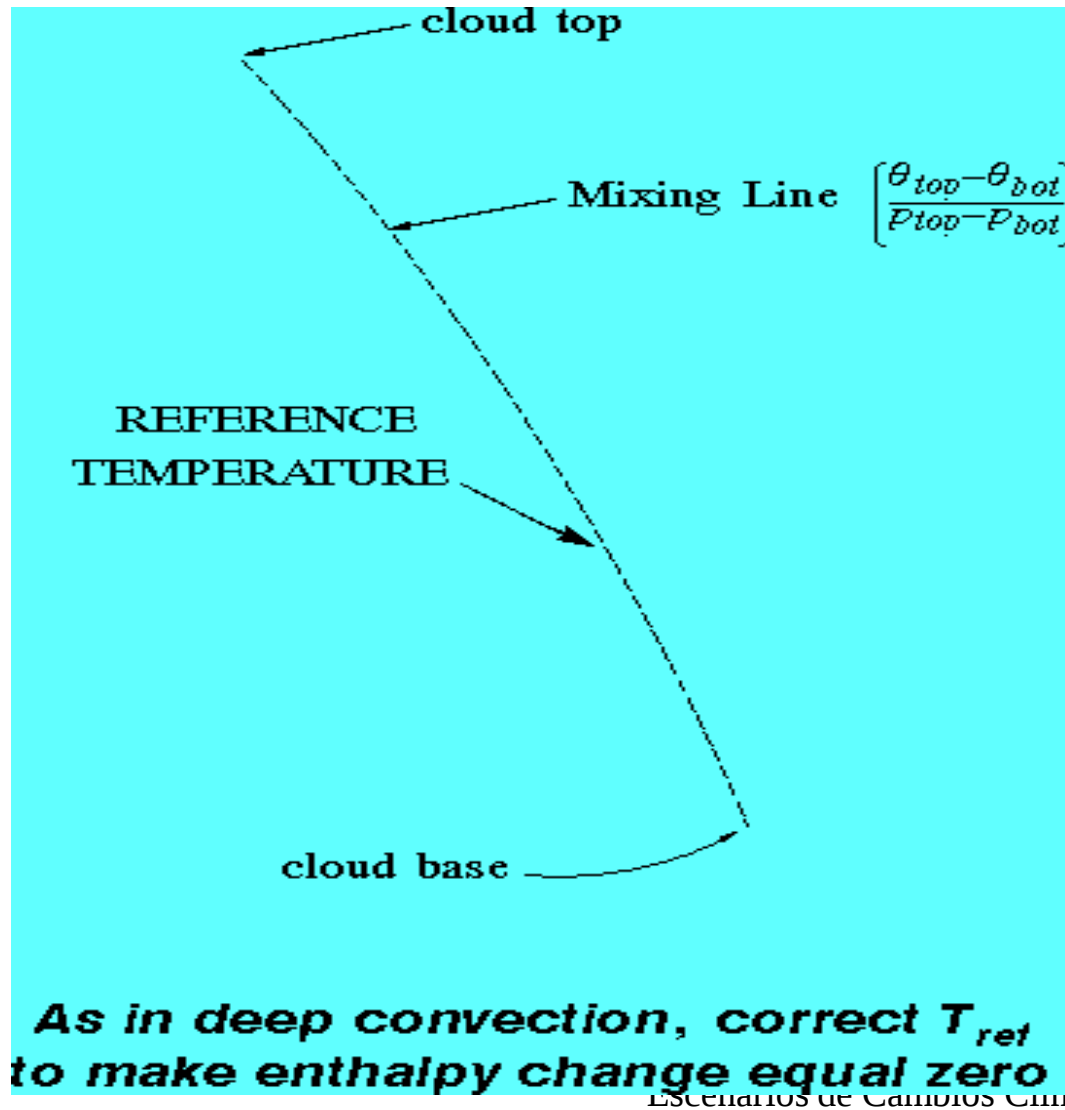
$$\tau' = \frac{\tau}{F(E)}$$

Fefi, fss



Shallow Convection

Temperature Reference Profile



- Applied to points where cloud depth is larger than 10hPa and smaller than 290hPa
- At least two layers
- swap points:
 - precipitation < 0
 - entropy change < 0

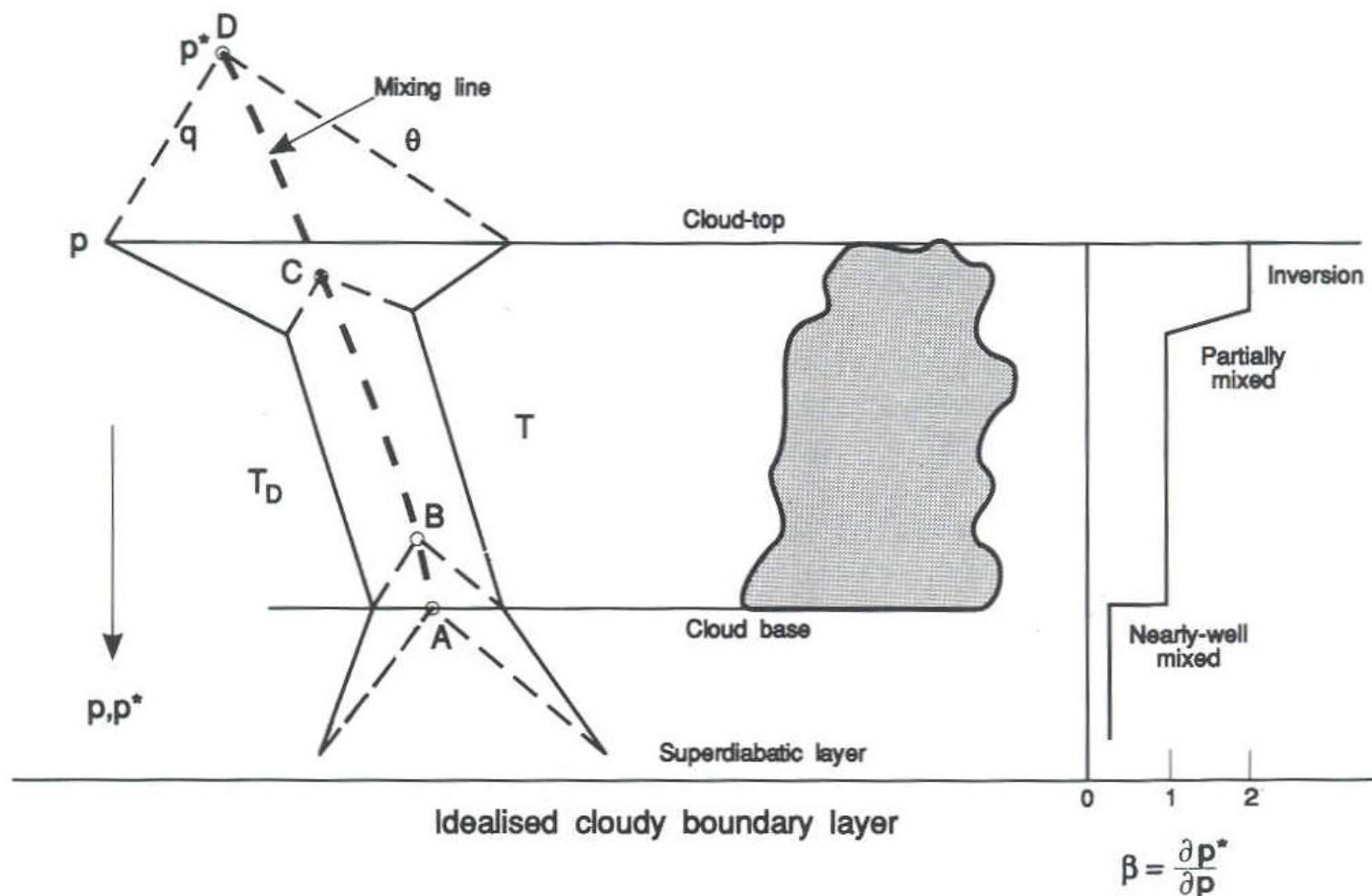
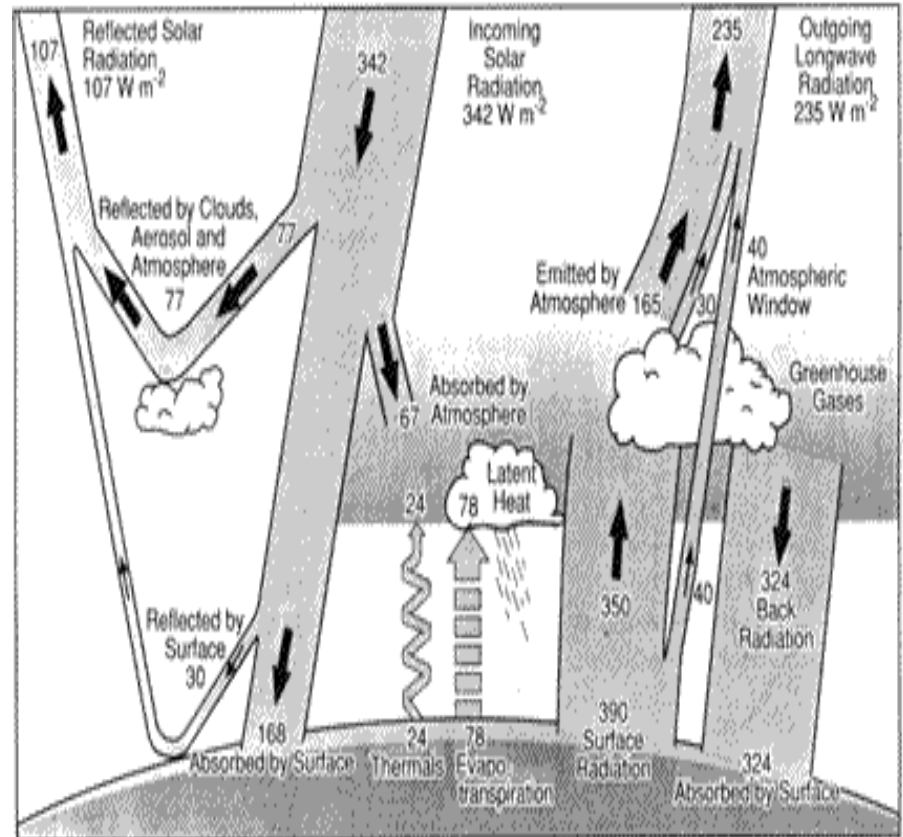


FIG. 9.1. Relationship between mixing line, temperature, and dewpoint, and a mixing parameter β for an idealized convective boundary layer. The light dashed lines are lines of constant potential temperature θ and mixing ratio q (from Betts 1986).

The need for modelling cloud microphysics

- *Presence of clouds change the radiation budget
- *Moisture and temperature adjustments in convectively stable points.
- *Precipitation production for model hydrology
- *Identification of precipitation types.

Excess from grid-point moisture excess in large scale models (large scale precipitation)



Cloud Formation

- Air rises and cools to saturation at LCL,
- Most effective Cloud Condensation Nuclei, CCN, are activated,
- Saturation vapor pressure decrease as parcel continues to rise and cool - the parcel becomes supersaturated,
- More CCN activate at the higher humidity.

Cloud Condensation Nuclei

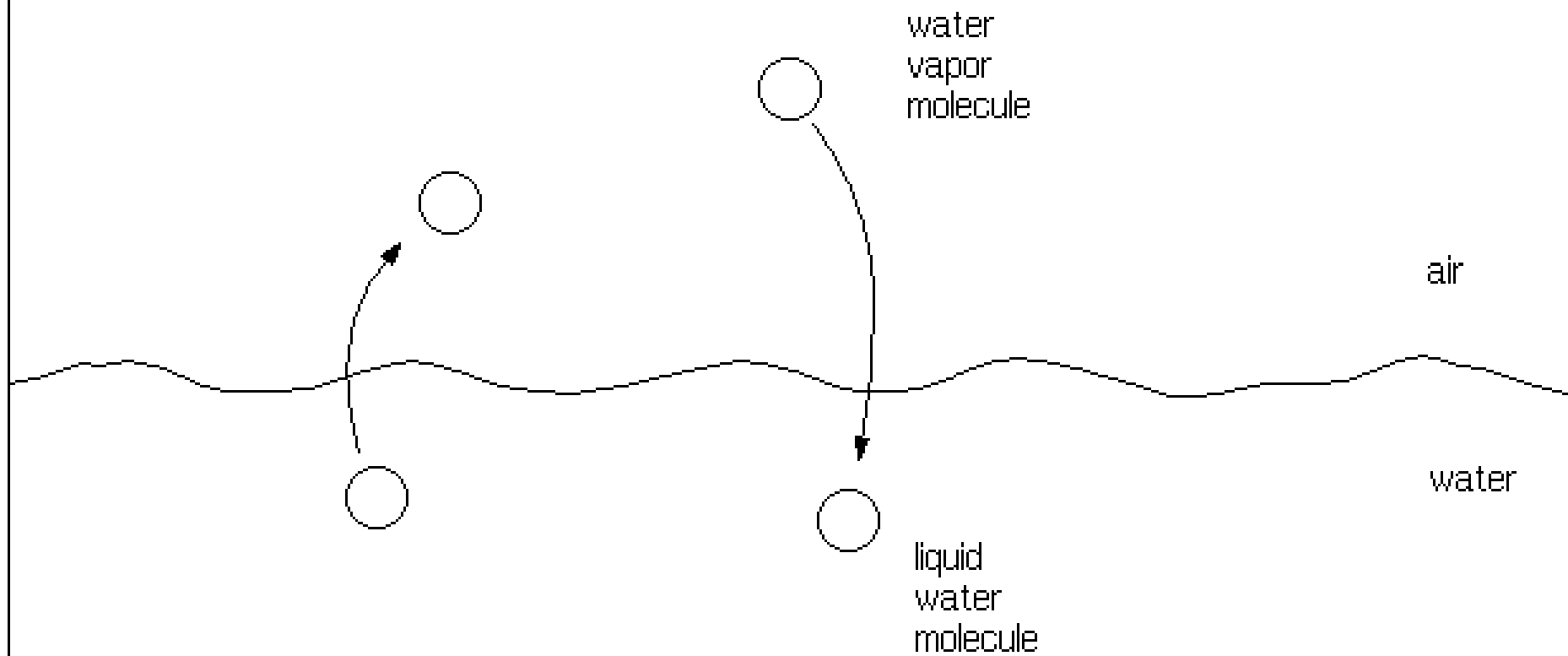
Particles suspended in the air which support the growth of cloud droplets or ice on their surface

soils, sand, dust, salt from sea spraying, volcanic debris, particles emitted from urban factories, natural aerosols

Cloud droplet Growth

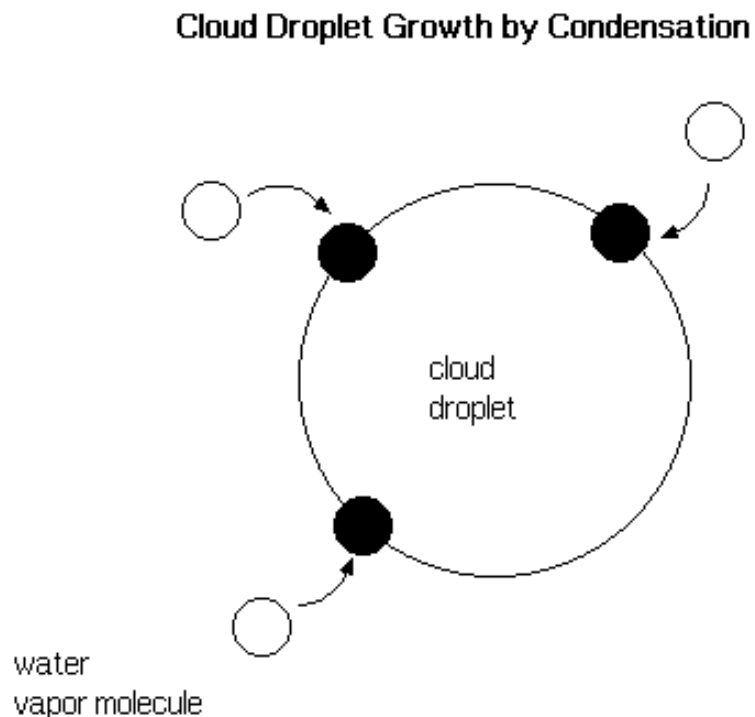
- Condensation - Diffusion
- Collision and Coalescence
- Melting from ice

Saturation



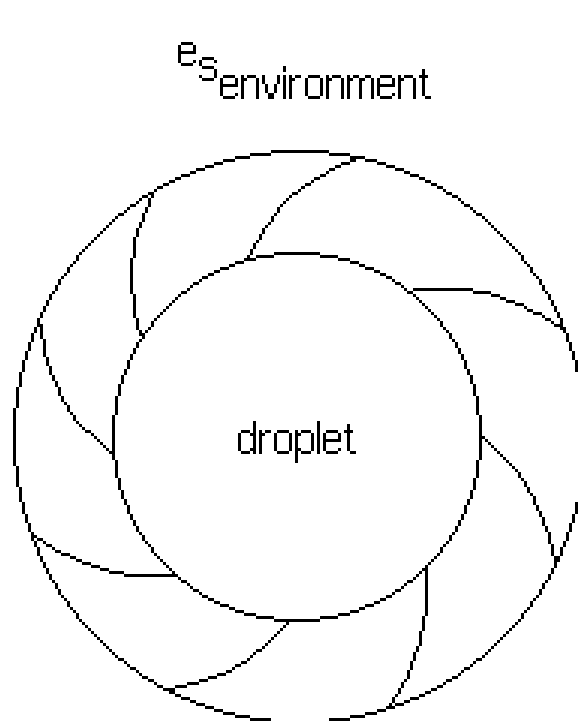
Cloud Droplet Growth by Condensation (Diffusion)

- Driven by the saturation vapor pressure difference
- Vapor is transported from higher to lower saturation vapor pressure.



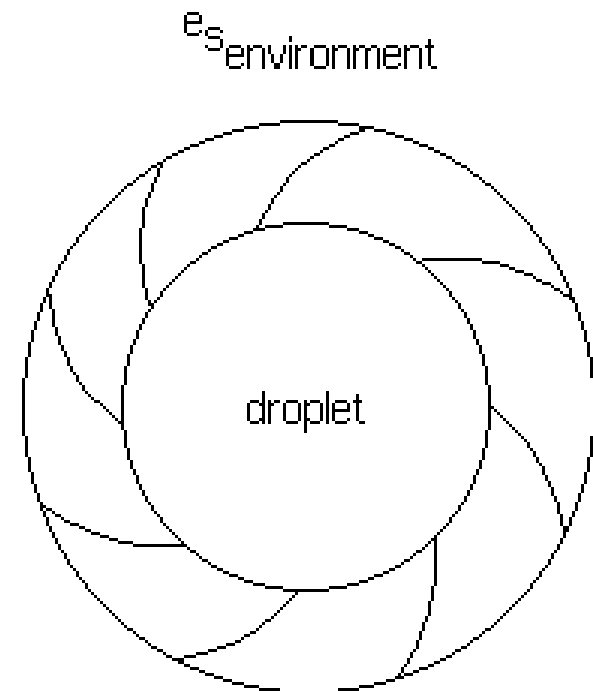
Comparisons between e_s of a Droplet and its Environment

the environment and the air surrounding the droplet
are assumed to be saturated



$$e_{s_{\text{env}}} > e_{s_{\text{droplet}}}$$

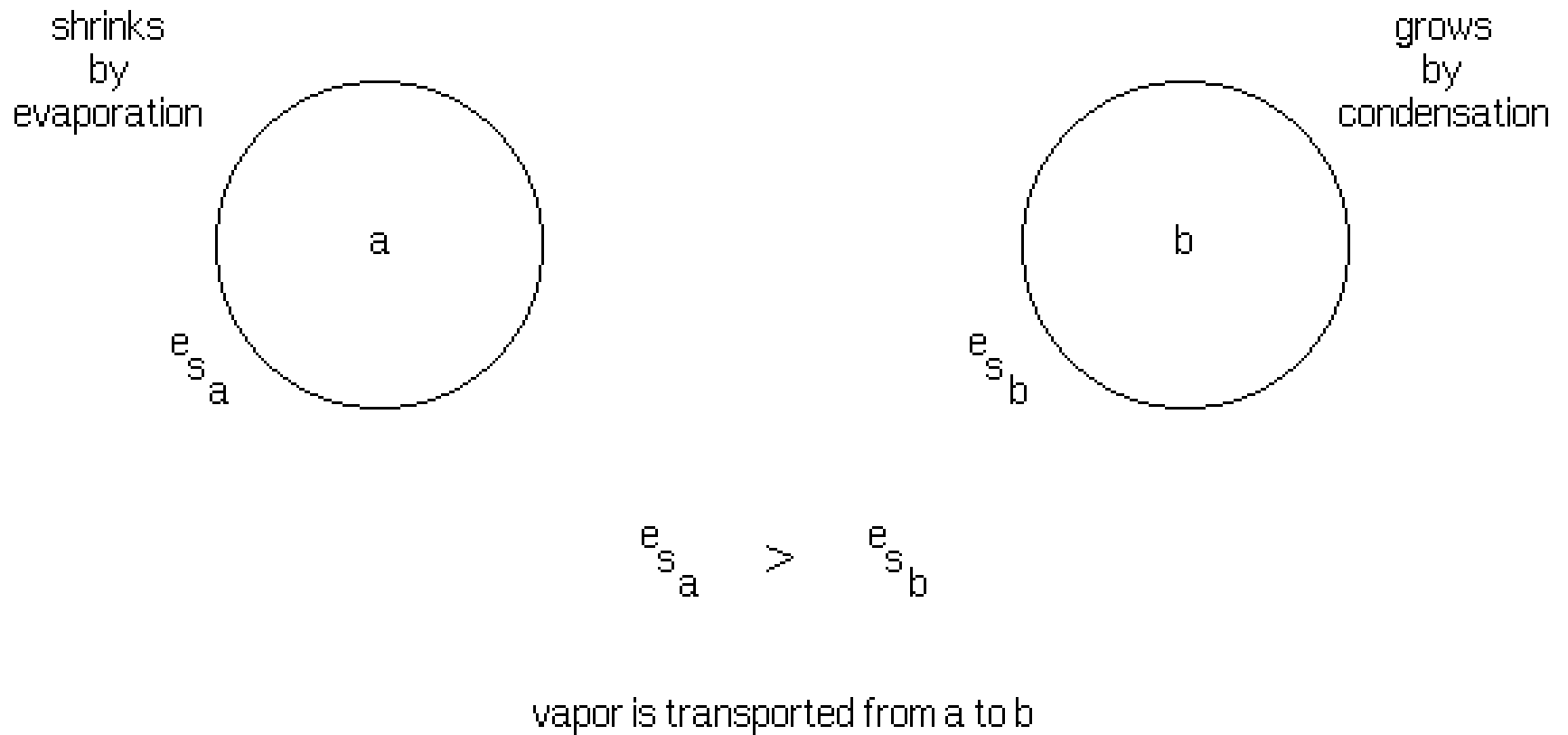
The droplet grows by
condensation



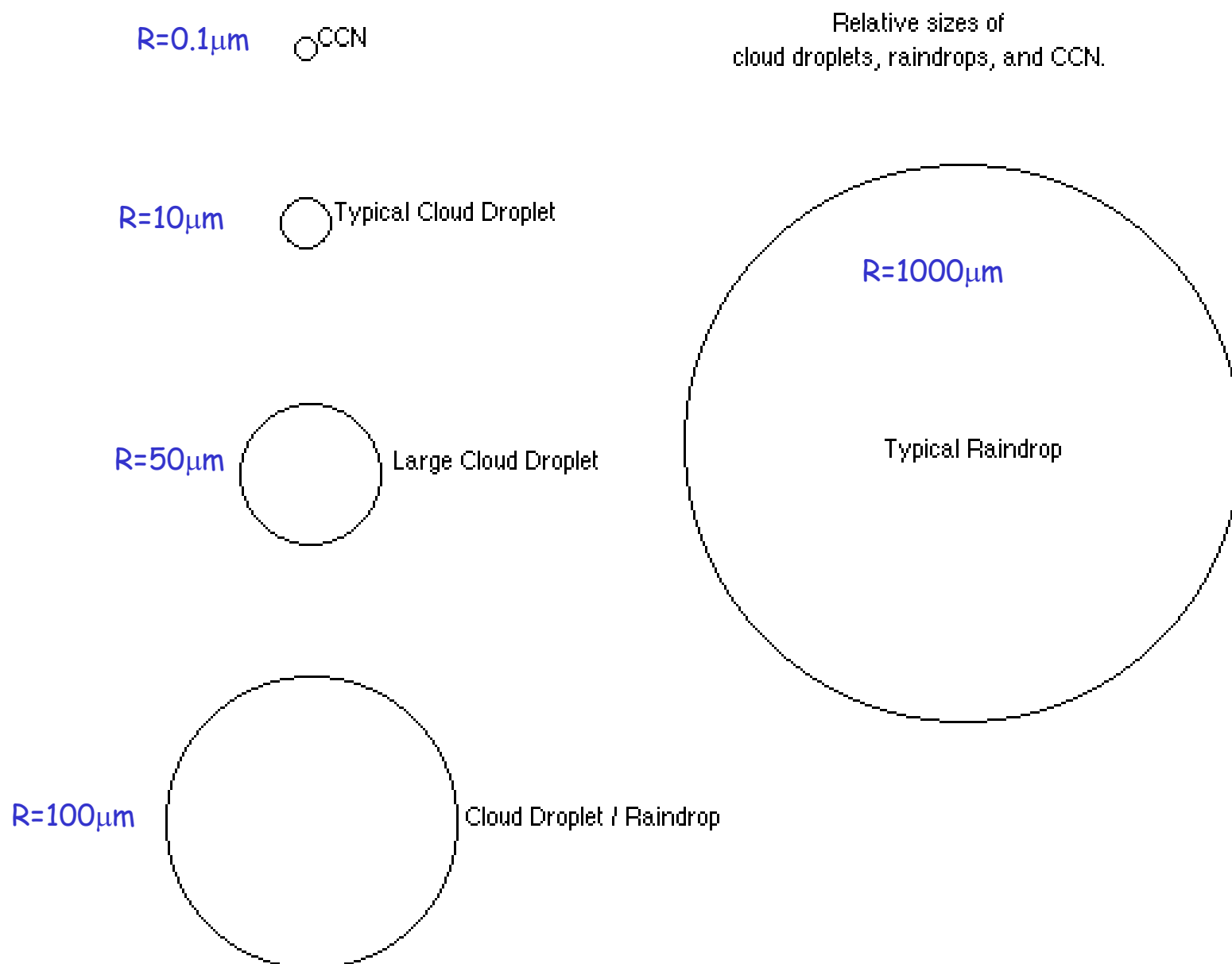
$$e_{s_{\text{env}}} < e_{s_{\text{droplet}}}$$

The droplet shrinks by
evaporation

Interaction of Droplets with different e_s



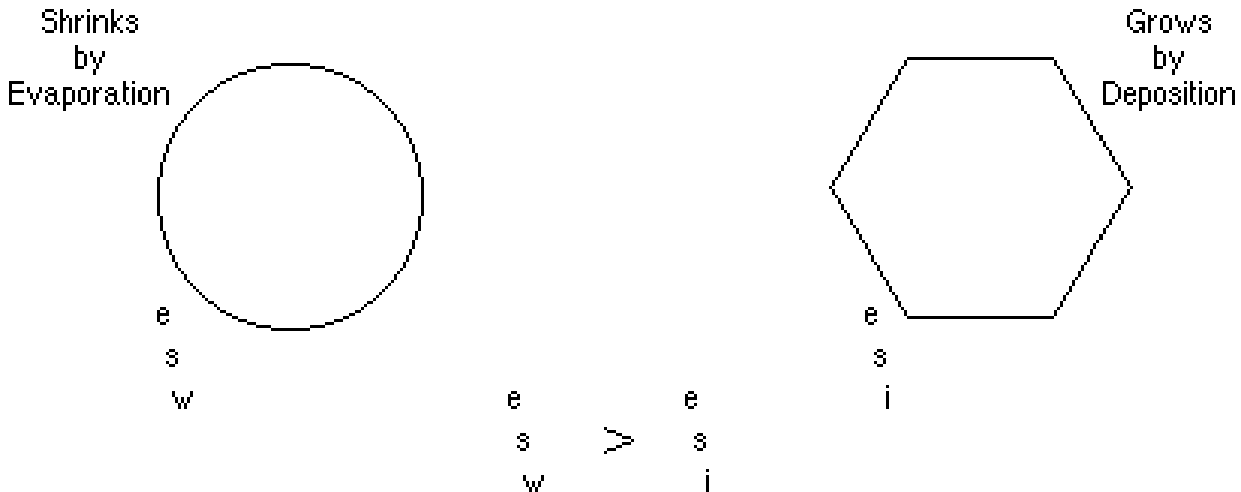
Comparative Sizes



Ice Crystal Growth

- Growth by **deposition**: Water vapor condenses onto an IN particle and freezes.
- Growth by **contact**: IN initiates an ice crystal at the moment it contacts a supercooled water droplet.
- Growth by **freezing**: IN floating within a supercooled droplet initiates the freezing.
- **Accretion** of cloud droplets: ice-phase particle has grown to a sufficient size to begin to fall and collect the supercooled droplet,
- **Aggregation**: ice particles coming together to form one main snowflake

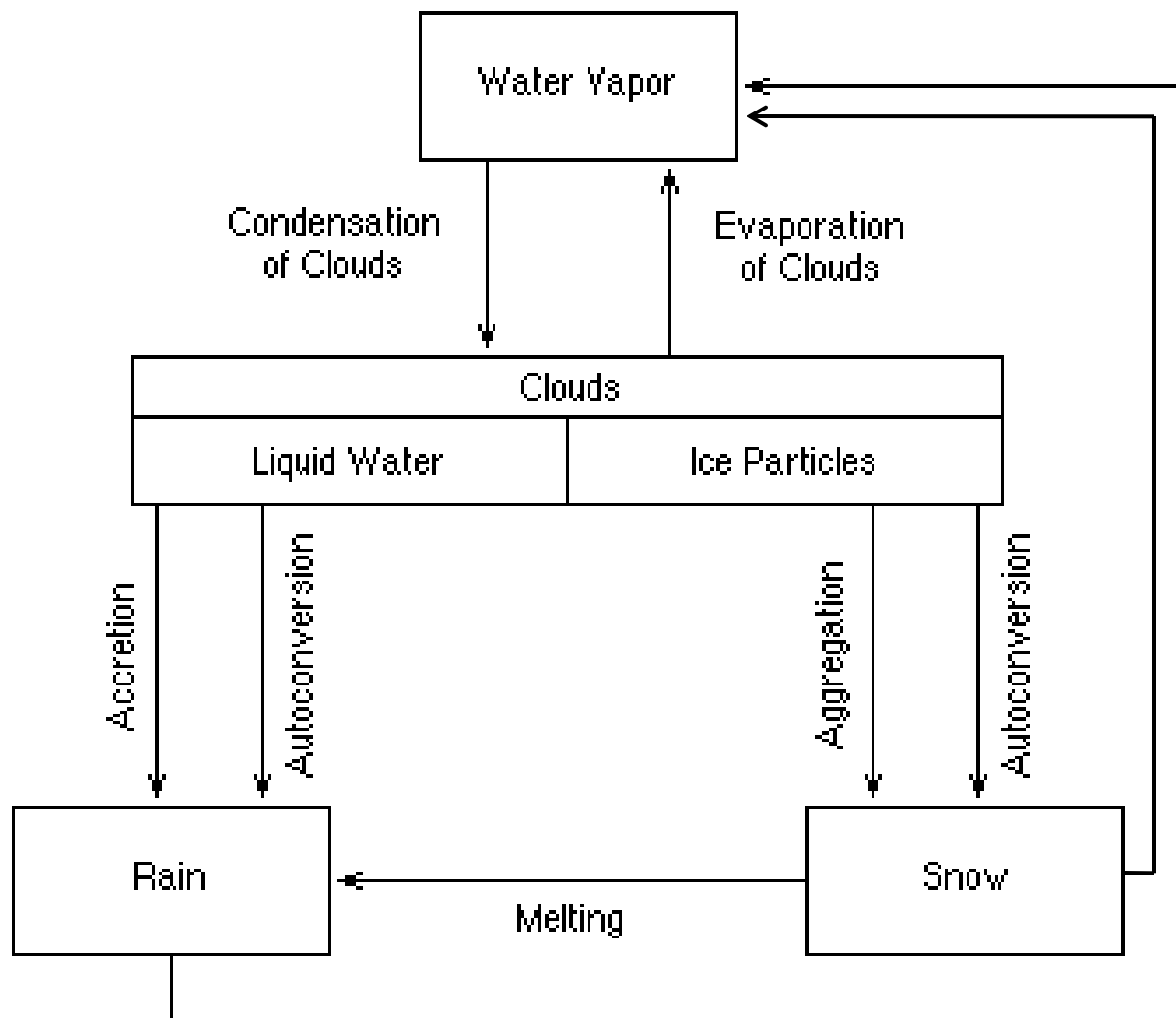
Interaction of Cloud Droplets and Ice Crystals



Vapor is transported from the droplet
to the ice crystal

conversions

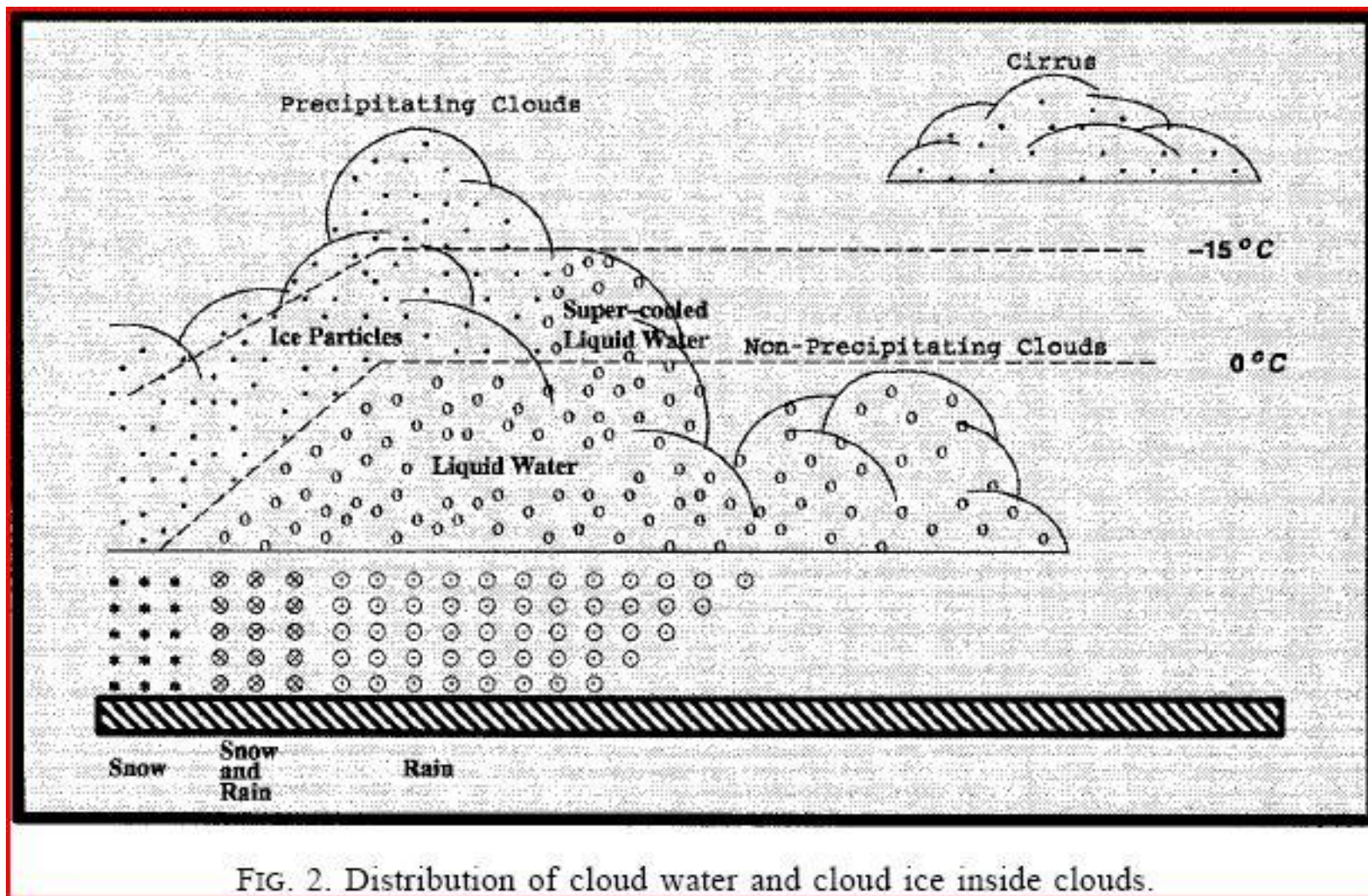
Water Vapor (Qv)	Snow (Qs)	Rain (Qr)	Cloud (Qc)	Ice (Qi)
condensation	sublimation	accretion	accretion	initiation
evaporation	accretion	autoconversion	autoconversion (into rain)	sublimation
ice initiation	autoconversion (from ice)	evaporation	condensation	deposition
deposition (onto ice crystals)	freezing (from rain)	freezing (to snow)	freezing (to ice)	melting (to cloud)
	melting (to rain)	melting (from snow)	melting (from ice)	freezing cloud
				accretion
				autoconversion



Zhao et al 1997

Zhao et al.

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Fração de cobertura de nuvens b

in Eq. (2.3). According to Sundqvist et al. (1989), cloud fraction b at a grid point can be estimated from relative humidity using the equation

$$b = 1 - \left(\frac{U_s - U}{U_s - U_{00}} \right)^{1/2} \quad (2.11)$$

for $U > U_{00}$ and $b = 0$ for $U < U_{00}$. Since both temperature

$$U_s = 1.0$$

$$U_{00} = 0.85$$

Precipitação P

Following Sundqvist et al. (1989), the autoconversion of cloud water to rain, P_{raut} , can be parameterized from the cloud water mixing ratio m and cloud fraction b ; that is,

$$P_{\text{raut}} = c_0 m \left\{ 1 - \exp \left[- \left(\frac{m}{m_r b} \right)^2 \right] \right\}, \quad (2.15)$$

where constants c_0 and m_r are $1.0 \times 10^{-4} \text{ s}^{-1}$ and 3.0×10^{-4} , respectively. The autoconversion of cloud ice to snow is simulated using the equation from Lin et al. (1983):

$$P_{\text{saat}} = a_1 (m - m_{i0}), \quad (2.16)$$

Taxa de precipitação por autoconversão da água da nuvem

Neve por autoconversão do gelo da nuvem

The collection of cloud liquid water by the falling rain is proportional to the cloud water mixing ratio m and the rain rate P_r , and can be expressed by

$$P_{\text{racw}} = C_r m P_r, \quad (2.18)$$

Taxa de precipitação por coleta de gotas de nuvem pela chuva em queda

$5.0 \times 10^{-4} \text{ m}^2 \text{ kg}^{-1} \text{ s}^{-1}$. Similarly, the aggregation process of ice particles by the falling snow is simulated by

$$P_{\text{saci}} = C_s m P_s, \quad (2.19)$$

where P_s is the precipitation rate of snow and C_s is the collection coefficient. Unlike C_r , C_s should be a function of temperature since the open structure of ice crystals at relatively warm temperatures increases the likelihood

Taxa de neve por
agregação cristais
de gelo com a neve
em queda

$$P_{\text{sm1}} = |C_{\text{sm}}(T - 273.15)^a P_s,$$

ter is given by (Zhao and Carr

$$P_{\text{sm2}} = C_{\text{ws}} P_{\text{sacw}},$$

is the collection rate of cloud li
g snow below the freezing level

$$P_{\text{sacw}} = C_r m P_s$$

Chuva por derretimento
da neve em queda

Chuva por derretimento
da neve formada por
coleta de água de nuvem

Neve por coleta de água
de nuvem pela neve

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Evaporation of rain as it falls through an unsaturated layer is calculated using the equation (Sundqvist 1988)

$$E_{rr} = k_e(U_{00} - U)(P_r)^\beta. \quad (2.24)$$

The sublimation of the falling snow is also computed using the equation

$$E_{rs} = [C_{rs1} + C_{rs2}(T - 273.15)] \frac{(U_{00} - U)}{U_{00}} P_s. \quad (2.25)$$

**Evaporação E
da chuva ou
neve, em
camadas
subsaturadas
subnuvem**

$$P_r^{(n)} = P_r^{(n-1)} + \frac{P_{sfc} - P_{top}}{g\eta_s} (P_{raut} + P_{racw} + P_{sacw} + P_{sml} + P_{sm2} - E_{rr})^{(n)} \Delta\eta \quad (2.26)$$

and

$$P_s^{(n)} = P_s^{(n-1)} + \frac{P_{sfc} - P_{top}}{g\eta_s} (P_{saut} + P_{sact} - P_{sml} - P_{sm2} - E_{rs})^{(n)} \Delta\eta, \quad (2.27)$$

Chuva e neve